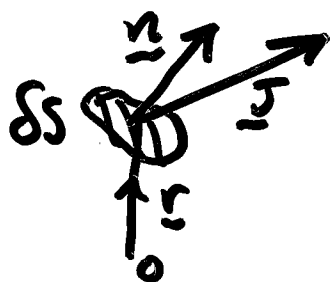


MATH2100 Lec. 25 (= MATH2011 Lec. 7)A conservation law:

Consider some substance flowing through space, with density $\rho(x, y, z, t)$ and flux vector field $\underline{J}(x, y, z, t)$:-



Amount of substance in δV at (x, y, z) at time $t = \rho(x, y, z, t) \delta V$

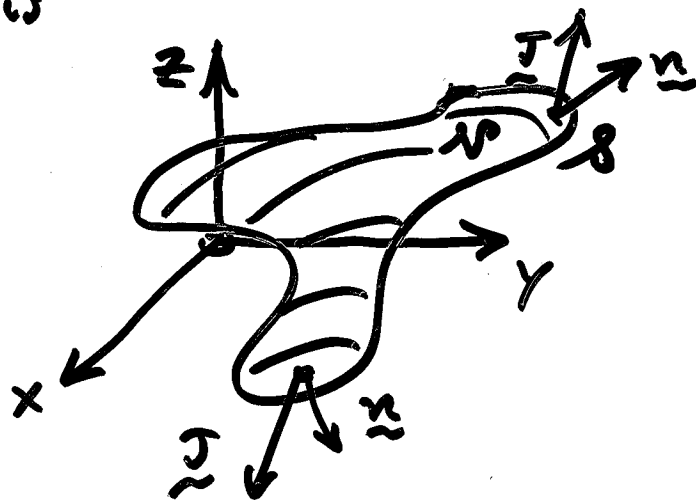


Amount of substance flowing per unit time over δS at (x, y, z)
 $= \underline{J}(x, y, z, t) \cdot \underline{n} \delta S$

(25.2)

If substance is conserved (not being created or destroyed), then for every closed surface S , and corresponding enclosed 3-D region V , we must have

$$\underbrace{\frac{d}{dt} \iiint_V \rho \, dv}_{\substack{\text{amount in } V \\ \text{at time } t \\ \text{rate at which} \\ \text{amount in } V \text{ is} \\ \text{increasing}}} = - \underbrace{\oint_S \underline{J} \cdot \underline{n} \, ds}_{\substack{\text{flux into } V \text{ over } S \\ \text{per unit time}}}$$



$$\Rightarrow \iiint_V \frac{\partial \rho}{\partial t} \, dv = - \iiint_V (\nabla \cdot \underline{J}) \, dv$$

using
Gauss'
Theorem

$$\Rightarrow \iiint_V \left(\frac{\partial \rho}{\partial t} + \nabla \cdot \underline{J} \right) dV = 0$$

Since S and corresponding V are arbitrary,
follows that at every point (x, y, z)
we have

$$\boxed{\frac{\partial \rho(x, y, z, t)}{\partial t} + \nabla \cdot \underline{J}(x, y, z, t) = 0} \quad (25.1)$$

Conservation of 'substance'

EX: 1) Moving fluid, velocity field $\underline{v}(x, y, z, t)$

If fluid is ^{homogeneous} incompressible, mass

density $\rho(x, y, z, t) = \rho_0$ (const.)

Since mass flux vector field is $\rho_0 \underline{v}$,
we have

$$\frac{\partial \rho_0}{\partial t} + \nabla \cdot (\rho_0 \underline{v}) = 0 \Rightarrow \rho_0 \nabla \cdot \underline{v} = 0$$

$$\Rightarrow \boxed{\nabla \cdot \underline{v} = 0}$$

Conservation of mass of incompressible fluid.

EX: 2)

25.4

Important example for rest of course:-

Conduction of heat in a uniform solid.

$u(x, y, z, t)$ = temperature at (x, y, z) at time t

ρ = mass/unit volume of solid

σ = specific heat of solid (= amount of heat energy to raise unit mass through one unit of temperature)

$\Rightarrow$ $\rho \sigma u(x, y, z, t)$ = heat energy density

Heat energy flux vector

$$\underline{J}(x, y, z, t) = -\kappa \underline{\nabla} u(x, y, z, t)$$

↑
thermal conductivity

[Heat flows from hotter to cooler regions, "down the temperature gradient"]

Conservation of heat energy:

$$\frac{\partial [\rho c u(x, y, z, t)]}{\partial t} + \nabla \cdot [-k \nabla u(x, y, z, t)] = 0$$

$$\Rightarrow \frac{\partial u(x, y, z, t)}{\partial t} - \underbrace{\left(\frac{k}{\rho c} \right)}_{"k=c^2"} \nabla \cdot [\nabla u(x, y, z, t)] = 0$$

$$\Rightarrow \boxed{\frac{\partial u(x, y, z, t)}{\partial t} = c^2 \nabla \cdot [\nabla u(x, y, z, t)]} \quad (25.2)$$

The heat (conduction) equation.

c^2 : thermal diffusivity

$$\begin{aligned} \text{Consider } \nabla \cdot [\nabla u] &= \nabla \cdot \left(\frac{\partial u}{\partial x} \underline{i} + \frac{\partial u}{\partial y} \underline{j} + \frac{\partial u}{\partial z} \underline{k} \right) \\ &= \left(\underline{i} \frac{\partial}{\partial x} + \underline{j} \frac{\partial}{\partial y} + \underline{k} \frac{\partial}{\partial z} \right) \cdot \left(\frac{\partial u}{\partial x} \underline{i} + \frac{\partial u}{\partial y} \underline{j} + \frac{\partial u}{\partial z} \underline{k} \right) \\ &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \end{aligned}$$

Write this as $\nabla^2 u$ ($= \text{del}^2 u$)

25.6

$$\nabla^2 = \nabla \cdot \nabla = \left(i \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right) \cdot \left(i \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right) \\ = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

∇ vector differential operator

$\nabla^2 = \nabla \cdot \nabla$ scalar differential operator

∇^2 called "the Laplacian"

So the heat equation (25.2) says

$$\frac{\partial u(x, y, z, t)}{\partial t} = c^2 \left[\frac{\partial^2 u(x, y, z, t)}{\partial x^2} + \frac{\partial^2 u(x, y, z, t)}{\partial y^2} + \frac{\partial^2 u(x, y, z, t)}{\partial z^2} \right] \\ = c^2 \nabla^2 u(x, y, z, t) \quad (25.3)$$

A partial differential equation (PDE)

4 independent variables x, y, z, t

1 dependent variable u

1st-order in t

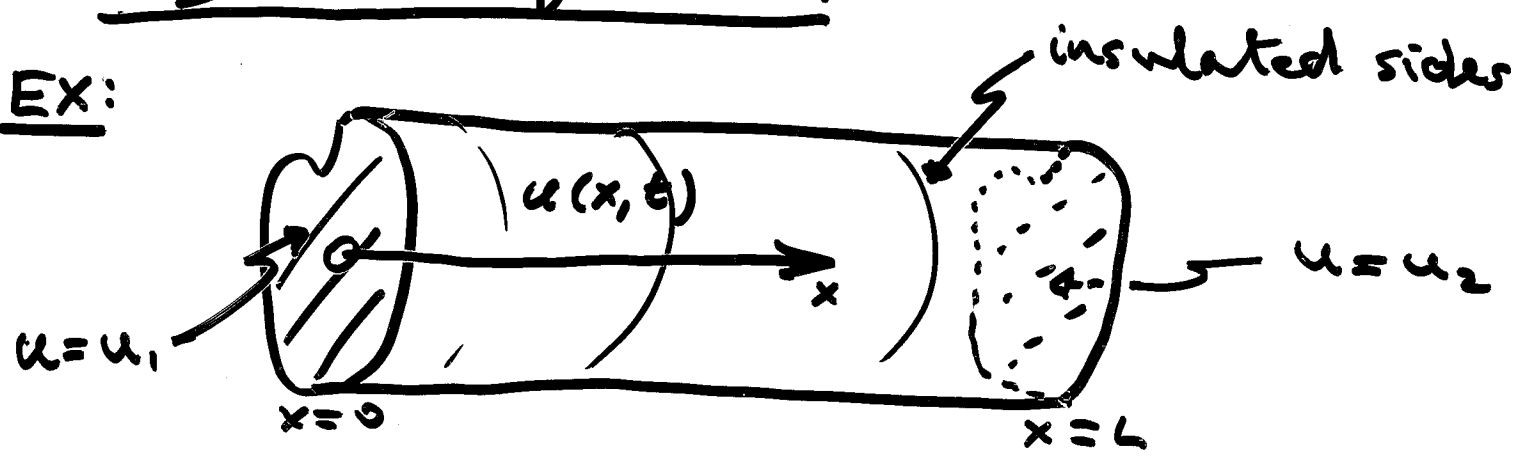
2nd order in x, y and z

Sometimes call (25.3) the 3-D heat equation
If no y or z dependence in problem,
get

$$\frac{\partial u(x,t)}{\partial t} = c^2 \frac{\partial^2 u(x,t)}{\partial x^2} \quad (25.4)$$

1-D heat equation.

EX:



iron bar, initially at constant temperature u_0 . Then one face held at temp. $u=u_1$, other at $u=u_2$.

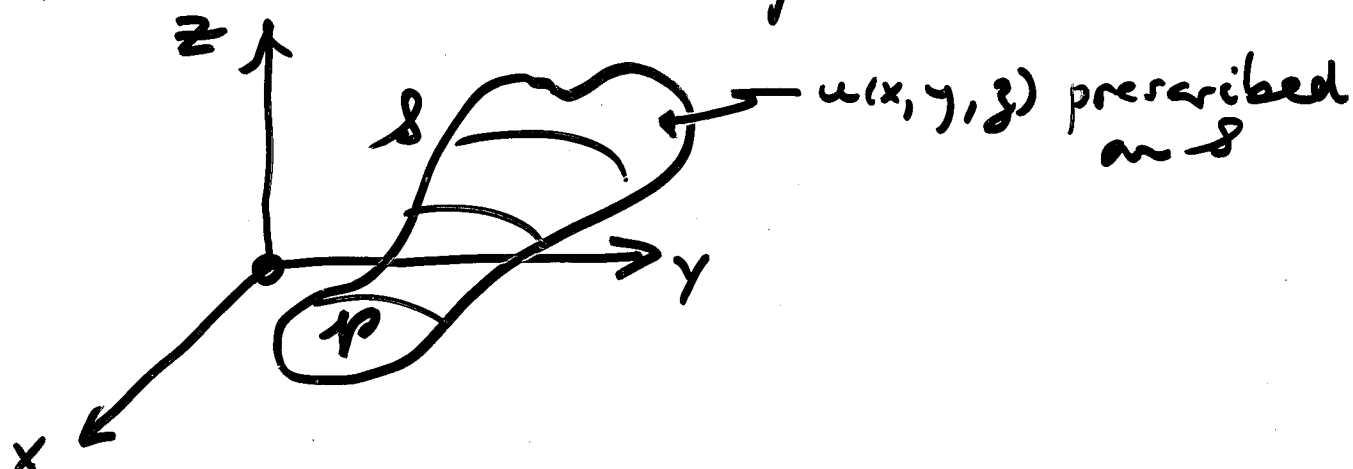
$u(x,t)$ satisfies (25.4) for $t > 0$ and $0 < x < L$.

Typical boundary/initial value problem:-

Find solution with IC $u(x,0) = u_0$
 $u(x,t)$ BC, $u(0,t) = u_1$, $u(L,t) = u_2$.

The heat equation governs the way temperature varies as heat flows from hotter to cooler regions, in a way consistent with conservation of heat energy.

Consider a 3-D region, with a prescribed, constant (in time) temperature distribution on the surface:—



After a long time, we expect temperature inside to approach a steady distribution $u(x, y, z)$ satisfying (from (25.3))

$$\frac{\partial u(x, y, z)}{\partial t} = c^2 \nabla^2 u(x, y, z)$$

$$\Rightarrow \nabla^2 u(x, y, z) = 0 \quad (25.5)$$

$$\frac{\partial^2 u(x, y, z)}{\partial x^2} + \frac{\partial^2 u(x, y, z)}{\partial y^2} + \frac{\partial^2 u(x, y, z)}{\partial z^2} = 0$$

Laplace's Equation

- describes equilibrium (steady-state) temperature distributions.

If no z -dependence in problem,

$$\frac{\partial^2 u(x, y)}{\partial x^2} + \frac{\partial^2 u(x, y)}{\partial y^2} = 0 \quad (25.6)$$

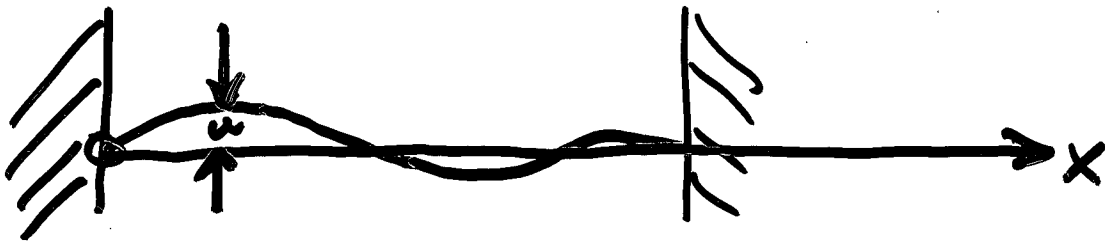
2-D Laplace Equation

PDE with 2 indep. variables x, y and one dependent variable u .

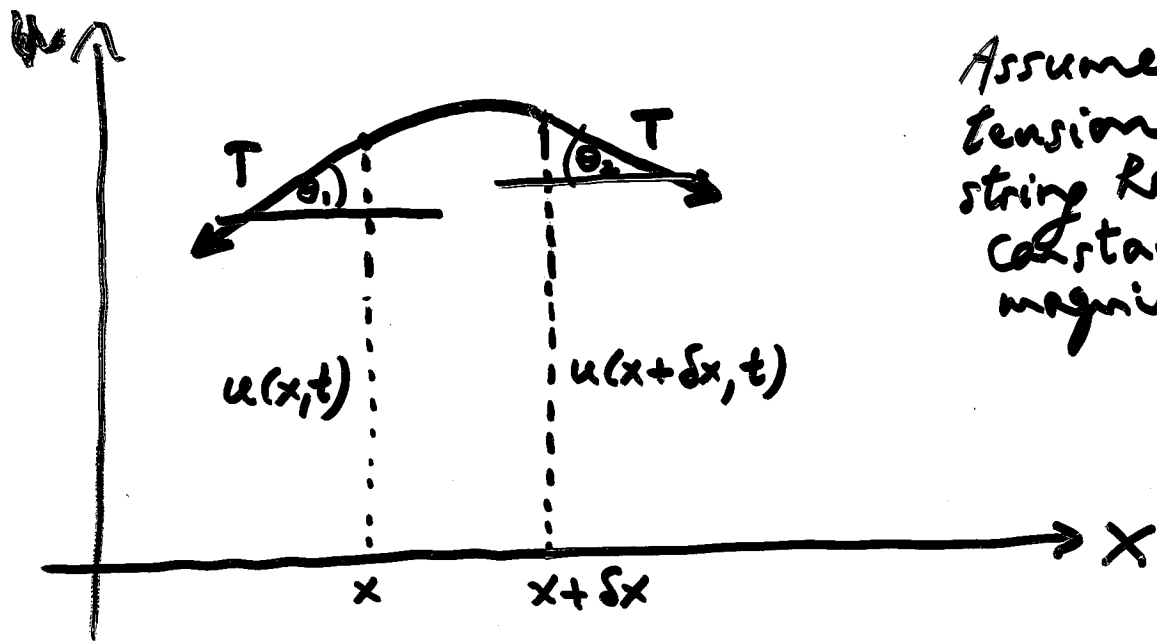
One more PDE with 2 indep. variables and one dep. variable for us to consider

Consider small displacements of stretched string, whose equilibrium position is along x -axis:-

25.10



Consider small segment:-



Assume
tension T in
string has
constant
magnitude

Force on segment in $+u$ -direction

$$\approx -T \sin \theta_1 - T \sin \theta_2$$

$$\approx -T \tan \theta_1 - T \tan \theta_2 \quad (|\theta_1|, |\theta_2| \ll 1)$$

$$= -T \frac{\partial u(x, t)}{\partial x} + T \frac{\partial u(x + \delta x, t)}{\partial x}$$

$$\approx -T \frac{\partial u(x, t)}{\partial x} + T \left[\frac{\partial u(x, t)}{\partial x} + \delta x \frac{\partial^2 u(x, t)}{\partial x^2} \right]$$

$$\approx T \frac{\partial^2 u(x, t)}{\partial x^2} \delta x$$

If mass/unit length of string is ρ , then
 mass $\times$ accel. of segment in u -direction
 is $(\rho \delta x) \frac{\partial^2 u(x, t)}{\partial t^2}$

Then, by Newton's 2nd Law,

$$T \frac{\partial^2 u(x, t)}{\partial x^2} \delta x \approx \rho \delta x \frac{\partial^2 u(x, t)}{\partial t^2}$$

$$\Rightarrow \boxed{\frac{\partial^2 u(x, t)}{\partial t^2} = c^2 \frac{\partial^2 u(x, t)}{\partial x^2} \quad (25.7)}$$

$$c^2 = \frac{T}{\rho}$$

1-D wave equation

[The 3-D wave equation is

$$\frac{\partial^2 u(x, y, z, t)}{\partial t^2} = c^2 \nabla^2 u(x, y, z, t)]$$

We will consider the 3 PDEs in one dependent, and two independent variables:-

The 1-D heat equation:

$$\frac{\partial u(x,t)}{\partial t} = c^2 \frac{\partial^2 u(x,t)}{\partial x^2}$$

Laplace's Equation in 2-D:

$$\frac{\partial^2 u(x,y)}{\partial x^2} + \frac{\partial^2 u(x,y)}{\partial y^2} = 0$$

The 1-D wave equation:

$$\frac{\partial^2 u(x,t)}{\partial t^2} = c^2 \frac{\partial^2 u(x,t)}{\partial x^2}$$

Summary:

- 1) Understand idea of a conservation law and know (25.1)
- 2) Know definition of Laplacian operator ∇^2
- 3) Know forms of heat eqn^s, wave eqn^s and Laplace's eqn^s in 1, 2, 3 dimensions

~~K pp 511, 512. § 11.1, 11.2~~
~~p 451~~

K pp. 407, 408, 411
 464-5
 § 12.1, 12.2