

MATH2011 — ANALYSIS OF PARTIAL DIFFERENTIAL EQUATIONS
Second Semester Examination, November, 2004 (continued)

Q 1 (a). Show that the Fourier Series corresponding to the function defined by

$$f(x) = \alpha, \quad -L < x \leq 0; \quad f(x) = \beta, \quad 0 < x \leq L;$$

and $f(x+2L) = f(x), \quad -\infty < x < \infty,$

where α and β are constants, is given by

$$\frac{1}{2}(\alpha + \beta) - \frac{2(\alpha - \beta)}{\pi} \sum_{m=1}^{\infty} \frac{1}{(2m-1)} \sin \left[\frac{(2m-1)\pi x}{L} \right]. \quad (*)$$

(17 marks)

$$f(x) = a_0 + \sum_{n=1}^{\infty} \left[a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right) \right] \quad (1)$$

$$a_0 = \frac{1}{2L} \int_{-L}^L f(x) dx = \frac{1}{2L} \int_{-L}^0 \alpha dx + \frac{1}{2L} \int_0^L \beta dx = \frac{1}{2}(\alpha + \beta) \quad (3)$$

$$\begin{aligned} a_n &= \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx \\ &= \frac{1}{L} \int_{-L}^0 \alpha \cos\left(\frac{n\pi x}{L}\right) dx + \frac{1}{L} \int_0^L \beta \cos\left(\frac{n\pi x}{L}\right) dx \\ &= \frac{1}{L} \left\{ \alpha \left[\frac{L}{n\pi} \sin\left(\frac{n\pi x}{L}\right) \right]_{-L}^0 + \beta \left[\frac{L}{n\pi} \sin\left(\frac{n\pi x}{L}\right) \right]_0^L \right\} \end{aligned} \quad (4)$$

$$= 0$$

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Q 1 (a). Working space only

$$\begin{aligned}
 b_n &= \frac{1}{2L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx \\
 &= \frac{1}{2L} \int_{-L}^0 \alpha \sin\left(\frac{n\pi x}{L}\right) dx + \frac{1}{2L} \int_0^L \beta \sin\left(\frac{n\pi x}{L}\right) dx \\
 &= \frac{1}{2L} \left\{ -\frac{\alpha L}{n\pi} \cos\left(\frac{n\pi x}{L}\right) \Big|_{-L}^0 - \frac{\beta L}{n\pi} \cos\left(\frac{n\pi x}{L}\right) \Big|_0^L \right\} \\
 &= -\frac{\alpha}{n\pi} \left[\underbrace{1 - \cos(-n\pi)}_{"1 - (-1)^n"} \right] - \frac{\beta}{n\pi} \left[\underbrace{\cos(n\pi) - 1}_{(-1)^n - 1} \right] \\
 &= -\frac{1}{n\pi} (\alpha - \beta) [1 - (-1)^n] \quad (7)
 \end{aligned}$$

$$b_1 = -\frac{1}{\pi} (\alpha - \beta) \cdot 2, \quad b_2 = 0, \quad b_3 = -\frac{1}{3\pi} (\alpha - \beta) \cdot 2, \dots$$

$$\begin{aligned}
 \text{Then } f(x) &= \frac{1}{2}(\alpha + \beta) - 2\left(\frac{\alpha - \beta}{\pi}\right) \left[1 \sin\left(\frac{\pi x}{L}\right) + \frac{1}{3} \sin\left(\frac{3\pi x}{L}\right) + \dots \right] \\
 &= \frac{1}{2}(\alpha + \beta) - 2\left(\frac{\alpha - \beta}{\pi}\right) \sum_{n=1}^{\infty} \frac{1}{(2n-1)} \sin\left[\frac{(2n-1)\pi x}{L}\right]. \quad (2)
 \end{aligned}$$

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Q 1 (b). Briefly explain the value taken by the series (*) at $x = 0$.

At $x=0$, series converges to $\frac{1}{2}(\alpha+\beta)$ (5 marks) ①
At the jump discontinuity, f has well defined limits from L and R , and so does f'

$$f(0_-) = \alpha, \quad f(0_+) = \beta \quad ①$$

$$f'(0_-) = 0, \quad f'(0_+) = 0. \quad ①$$

Then we know series will converge to average value $\frac{1}{2}(\alpha+\beta)$ across jump ②

Q 1 (c). By considering the value to which the series (*) must converge at $x = L/2$ deduce that

$$1 - \frac{1}{3} + \frac{1}{5} - \dots = \frac{\pi}{4}.$$

(6 marks)

We know series will converge to $f(\frac{1}{2}L) = \beta$ as function is cont. & diffble there. ②

$$\therefore \beta = \frac{1}{2}(\alpha+\beta) - 2\left(\frac{\alpha-\beta}{\pi}\right) \sum_{m=1}^{\infty} \left(\frac{1}{2m-1}\right) \sin\left[\frac{(2m-1)\pi}{2}\right] \quad ②$$

$$\frac{1}{2}(\beta-\alpha) = 2\left(\frac{\beta-\alpha}{\pi}\right) \left[1 - \frac{1}{3} + \frac{1}{5} - \dots\right]$$

Must hold whether or not $\alpha \neq \beta$. \therefore

$$\frac{1}{2} = \frac{2}{\pi} \left[1 - \frac{1}{3} + \frac{1}{5} - \dots\right]$$

$$\Rightarrow 1 - \frac{1}{3} + \frac{1}{5} - \dots = \frac{\pi}{4} \quad ②$$

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Q 1 (d). Describe and explain in a few words what happens to the series (*) when (a) $\alpha = \beta$ and (b) $\alpha = -\beta$.

(5 marks)

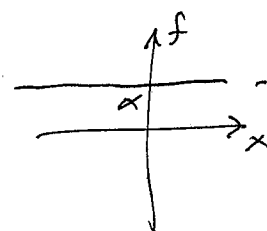
(a) When $\alpha = \beta$.

$f(x) = \alpha$ (const.) — even fn

series collapse to value α .

— even function with
~~trivial~~ cosine series

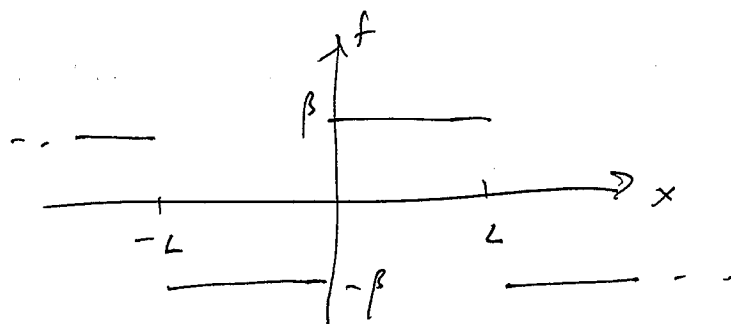
(2)



(b) when $\alpha = -\beta$.

$f(x)$ is odd fn

with sine series:



$$f(x) = \frac{4\beta}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1)} \sin\left[\frac{(2n-1)\pi x}{L}\right]$$

(3)

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Q 2 (a). You are given (no need to check!) that the function $G(x - y, t)$ defined by

$$G(x - y, t) = \frac{1}{\sqrt{4\pi c^2 t}} e^{-(x-y)^2/(4c^2 t)}$$

satisfies

$$G_t(x - y, t) = c^2 G_{xx}(x - y, t), \quad -\infty < x < \infty, \quad t > 0.$$

Show that $u(x, t)$ defined by

$$u(x, t) = \int_{-\infty}^{\infty} G(x - y, t) f(y) dy$$

satisfies the 1-dimensional Heat Equation for $-\infty < x < \infty$ and $t > 0$, and also the initial condition

$$\lim_{t \rightarrow 0^+} u(x, t) = f(x), \quad -\infty < x < \infty.$$

(9 marks)

$$u(x, t) = \int_{-\infty}^{\infty} G(x - y, t) f(y) dy$$

$$u_t(x, t) = \int_{-\infty}^{\infty} G_t(x - y, t) f(y) dy$$

$$u_{xx}(x, t) = \int_{-\infty}^{\infty} G_{xx}(x - y, t) f(y) dy$$

$$\therefore u_t - c^2 u_{xx} = \int_{-\infty}^{\infty} [G_t(x - y, t) - c^2 G_{xx}(x - y, t)] f(y) dy$$

$$= 0.$$

(3)

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Q 2 (a). Working space only

$$u(x, t) = \frac{1}{\sqrt{4\pi c^2 t}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4c^2 t} f(y) dy$$

Put $v = \frac{x-y}{\sqrt{4c^2 t}} \Leftrightarrow y = x - \sqrt{4c^2 t} v$
 $dy = -\sqrt{4c^2 t} dv$

$$y = -\infty \Leftrightarrow v = \infty$$

$$y = +\infty \Leftrightarrow v = -\infty$$

$$\therefore u(x, t) = \frac{1}{\sqrt{4\pi c^2 t}} \int_{-\infty}^{\infty} e^{-v^2} f(x - \sqrt{4c^2 t} v) (-dv) \sqrt{4c^2 t}$$

$$= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-v^2} f(x - \sqrt{4c^2 t} v) dv \quad (3)$$

As $t \rightarrow 0_+$

$$u(x, t) \Rightarrow \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-v^2} f(x) dv$$

$$= f(x) \cdot \underbrace{\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-v^2} dv}_1$$

$$= f(x)$$

(3)

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Q 2 (b). A very long cylindrical iron bar, with unknown thermal diffusivity c^2 , lies along the positive x -axis. The sides of the bar and the end at $x = 0$ are thermally insulated, and the temperature u inside the bar is a function only of x and of time t . The temperature distribution in the bar at $t = 0$ is given by

$$u(x, 0) = F(x) = \begin{cases} u_0 \text{ (const.)} & , \quad 0 < x < L \\ 0 & , \quad L < x < \infty \end{cases}$$

Deduce that

$$u(x, t) = \frac{1}{2} u_0 \left\{ \operatorname{erf} \left(\frac{x+L}{\sqrt{4c^2 t}} \right) - \operatorname{erf} \left(\frac{x-L}{\sqrt{4c^2 t}} \right) \right\}$$

for $x > 0, t > 0$, where

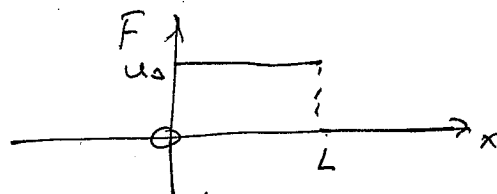
$$\operatorname{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-v^2} dv.$$

(18 marks)

We have to solve

$$u_t(x, t) = c^2 u_{xx}(x, t) \quad x > 0, t > 0$$

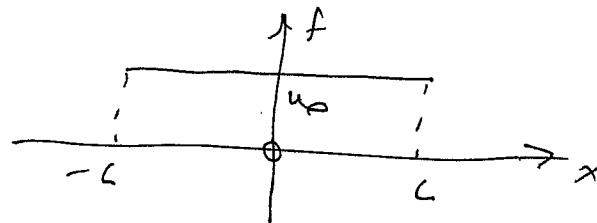
$$\left. \begin{array}{l} \text{BC} \quad u_x(0, t) = 0 \\ \text{IC} \quad u(x, 0) = F(x) \end{array} \right\} \quad (3)$$



Extend to problem on whole x -axis with even initial data, so BC is satisfied automatically: -

$$\text{So consider } u_t(x, t) = c^2 u_{xx}(x, t) \quad -\infty < x < \infty, t > 0$$

$$\text{IC } u(x, 0) = f(x)$$



$$\underline{\underline{\text{Sol}}}: \quad u(x, t) = \frac{1}{\sqrt{4\pi c^2 t}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4c^2 t} f(y) dy \quad (1)$$

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Q 2 (b). Working space only

$$= \frac{u_0}{\sqrt{4\pi c^2 t}} \int_{-L}^L e^{-(x-y)^2/4c^2 t} dy$$

$$v = \frac{x-y}{\sqrt{4c^2 t}} \Leftrightarrow y = x - \sqrt{4c^2 t} v$$

$$dy = -\sqrt{4c^2 t} dv \quad (3)$$

$$y = \pm L \Leftrightarrow v = \frac{x \mp L}{\sqrt{4c^2 t}}$$

$$\int u(x,t) = \frac{u_0}{\sqrt{4\pi c^2 t}} \int_{\frac{x-L}{\sqrt{4c^2 t}}}^{\frac{x+L}{\sqrt{4c^2 t}}} e^{-v^2} (-\sqrt{4c^2 t}) dv \quad (1)$$

$$= \frac{u_0}{\sqrt{\pi}} \int_{\frac{x-L}{\sqrt{4c^2 t}}}^{\frac{x+L}{\sqrt{4c^2 t}}} e^{-v^2} dv \quad (2)$$

$$= \frac{1}{2} u_0 \left\{ \frac{2}{\sqrt{\pi}} \int_0^{\frac{x+L}{\sqrt{4c^2 t}}} e^{-v^2} dv - \frac{2}{\sqrt{\pi}} \int_0^{\frac{x-L}{\sqrt{4c^2 t}}} e^{-v^2} dv \right\} \quad (3)$$

$$= \frac{1}{2} u_0 \left\{ \operatorname{erf} \left(\frac{x+L}{\sqrt{4c^2 t}} \right) - \operatorname{erf} \left(\frac{x-L}{\sqrt{4c^2 t}} \right) \right\} \quad (2)$$

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Q 2 (c). Given that $\text{erf}(0.5) \approx 0.5$, deduce that the thermal diffusivity c^2 is given approximately by

$$c^2 = L^2/t_1,$$

where t_1 is the time taken for the temperature at the face $x = 0$ to reach the value $u_0/2$.

(5 marks)

~~At~~ $x=0$

$$\begin{aligned} u(0, t) &= \frac{1}{2} u_0 \left\{ \text{erf} \left(\frac{L}{\sqrt{\pi c^2 t}} \right) - \text{erf} \left(\frac{-L}{\sqrt{\pi c^2 t}} \right) \right\} \\ &= u_0 \text{erf} \left(\frac{L}{\sqrt{\pi c^2 t}} \right) \quad \text{as } \text{erf}(-z) = -\text{erf}(z) \end{aligned}$$

(2)

Then

$$u(0, t_1) = \frac{1}{2} u_0 = u_0 \text{erf} \left(\frac{L}{\sqrt{\pi c^2 t_1}} \right)$$

$$\Leftrightarrow \text{erf} \left(\frac{L}{\sqrt{\pi c^2 t_1}} \right) = \frac{1}{2}$$

$$\Leftrightarrow \frac{L}{\sqrt{\pi c^2 t_1}} \approx \frac{1}{2}$$

$$\Leftrightarrow c^2 \approx \frac{L^2}{t_1}$$

(3)

