

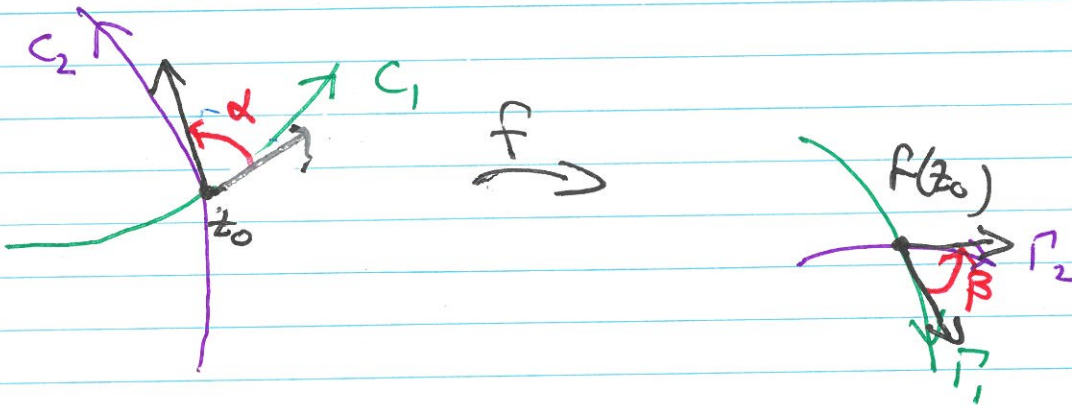
LECTURE 24§112 (8 Ed §101) Conformal Maps

- * $f: z \mapsto w$;
- * f analytic;
- * $f'(z_0) \neq 0$.

Then locally (near z_0) f preserves

- * angle;
- * orientation;
- * shape.

f is called conformal at z_0 .

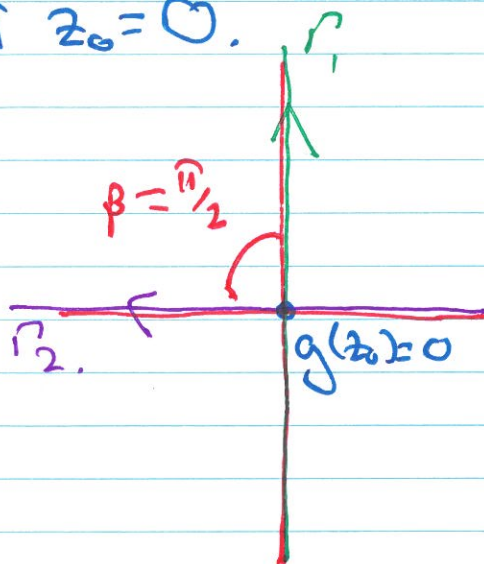
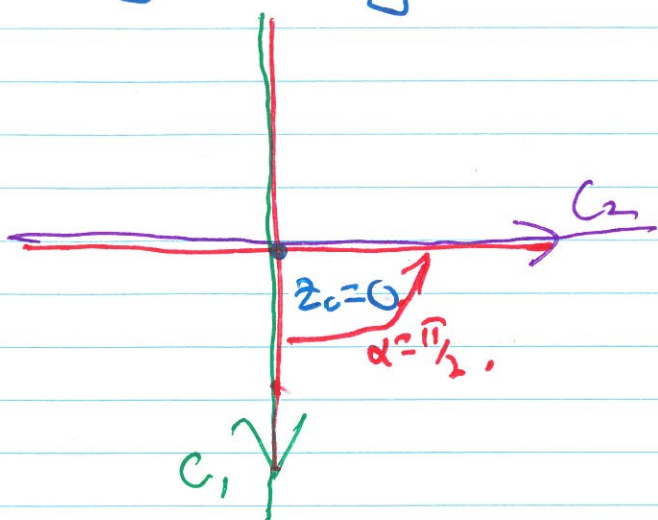


$$\Gamma_1 = f(C_1), \quad \Gamma_2 = f(C_2).$$

Conformality $\Rightarrow \beta = \alpha$

Roughly: $w = f(z), z = z(t)$:
 $w' = f'(z(t)) \cdot z'(t)$.
 $\Rightarrow \arg w' = \arg f' + \arg z'$

E.g: $g: z \mapsto -z$ at $z_0 = 0$.



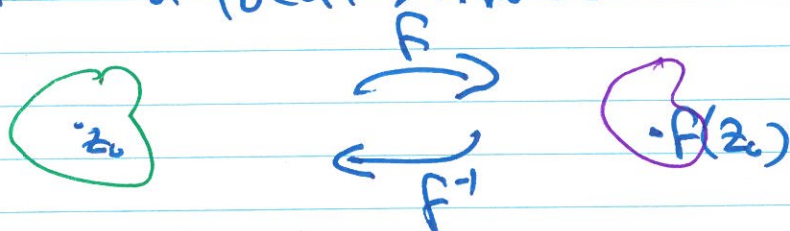
Analytic $f: D \rightarrow \mathbb{C}$ with $f'(z_0) = 0$: z_0 is a critical value of f . Angle will not be preserved.

Can show: angle will be multiplied by m , where m is the smallest integer s.t.

$$f^{(m)}(z_0) \neq 0.$$

RMK: If orientation is not necessarily preserved, but angle magnitude is, the map is called isogonal, e.g., $z \mapsto \bar{z}$.

Conformality \Rightarrow locally 1-1 & onto, i.e., f has a local inverse.



Inverse f^{-1} : thm (MATH2400/2401) in \mathbb{R}^2

Says: $f: (x,y) \mapsto (u,v)$ sufficiently smooth is invertible if

$$\det(\underline{J}_f) = \begin{vmatrix} u_x & u_y \\ v_x & v_y \end{vmatrix} \neq 0.$$

Jacobian

In our case: f is analytic $\Rightarrow u_x, u_y, v_x, v_y$ are all cts in a nbhd of $z_0 = x_0 + iy_0$,

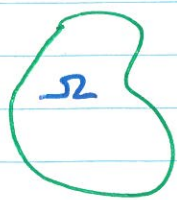
$$\det(\underline{J}_f) = \begin{matrix} u_x v_y - u_y v_x \\ \parallel \quad \parallel \\ u_x \quad -v_x \end{matrix}$$

$$= u_x^2 + v_x^2$$

$$= |u_x + iv_x|^2$$

$$= |f'|^2 \neq 0 \text{ at } z_0. \quad \square$$

Harmonic f's in \mathbb{R}^n



$\Omega \subseteq \mathbb{R}^n$; look $U : \Omega \rightarrow \mathbb{R}$

s.t. $\Delta U = 0$

Laplacian or Laplace operator.
(sometimes ∇^2)

where $\Delta U = \sum_{j=1}^n U_{jj} = \sum_{j=1}^n \frac{\partial^2 U}{\partial x_j^2}$
del or nabla

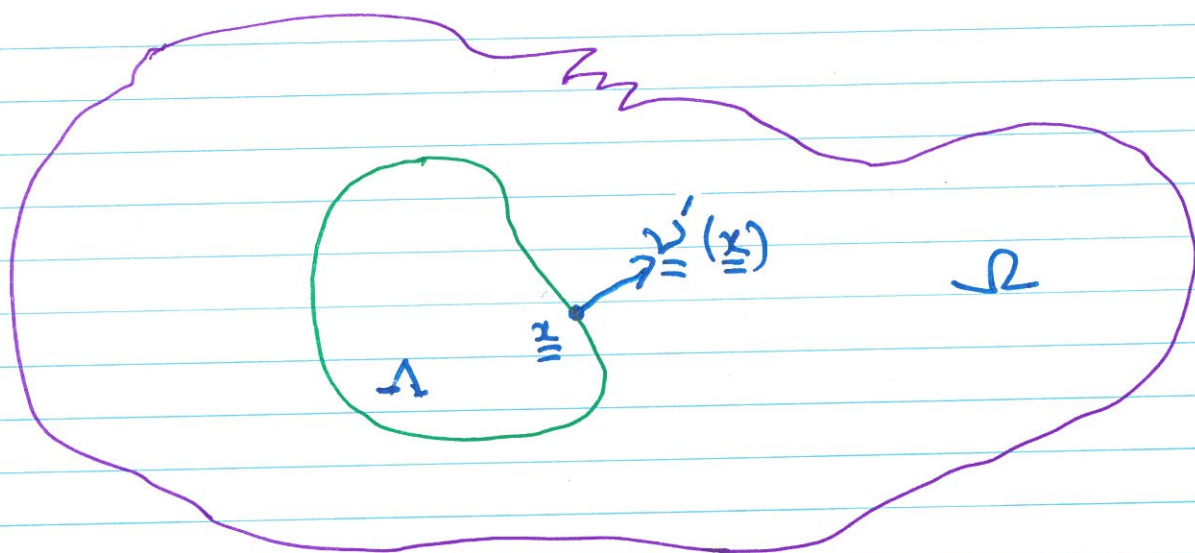
in \mathbb{R}^2 : $\Delta U = U_{xx} + U_{yy}$;

in \mathbb{R}^3 : $\Delta U = U_{xx} + U_{yy} + U_{zz}$.

Models many physical situations in
 "steady state".

Motivation: $\Omega \subseteq \mathbb{R}^2$ or \mathbb{R}^3 .

Λ : "sufficiently smooth" subdomain of Ω
 with an exterior normal $\underline{\underline{v}}(\underline{\underline{x}})$ on $\partial\Lambda$, $\underline{\underline{v}}$
 external unit normal $\underline{\underline{v}}'(\underline{\underline{x}})$.
 " $\underline{\underline{v}}$ " denotes a vector quantity.



" $\underline{\underline{u}}$ density of something" in equilibrium in Ω
 F = Flux density of " $\underline{\underline{u}}$ in equilibrium" in Ω .

$$\int_{\partial\Lambda} \underline{\underline{F}} \cdot \underline{\underline{v}}' dS = 0 \quad dS = \text{surface measure on } \partial\Lambda$$

$$\text{Gauss} \Rightarrow \int_{\Lambda} \text{div } \underline{\underline{F}} d\underline{\underline{x}} = 0$$

$$\begin{aligned} d\underline{\underline{x}} &= dx dy \quad \text{in 2D} \\ &= dx dy dz \quad \text{in 3D.} \end{aligned}$$

Since Λ is (essentially) arbitrary, these must hold:

$$\text{div } \underline{\underline{F}} = 0 \text{ in } \mathcal{R},$$

i.e., $\sum_{j=1}^n \partial_j F_j = 0 \text{ in } \mathcal{R}. \quad (*)$

In many physical situations,

$$\underline{\underline{F}} = c \nabla \bar{U}, \text{ with } c \text{ typically } -ve$$

$$(*) \Rightarrow c (\text{div } \nabla \bar{U}) = 0$$

$$\text{i.e., } \Delta \bar{U} = 0.$$

Can also study: $\frac{\partial \bar{U}}{\partial t} = a \Delta \bar{U} \quad (*)$

- If \bar{U} is the concentration of a chemical, $(*)$ is Fick's law of chemical diffusion;
- If \bar{U} is temperature, $(*)$ is Fourier's law of heat conduction;
- If \bar{U} is electric potential, $(*)$ is Ohm's law of electrical conduction.