

- Q1. (a) Derive Hamilton's equations of motion from Lagrange's equations of motion.  
(3 marks)

Lagrange's equations of motion are

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0 \quad \text{where } L(q_i, \dot{q}_i, t).$$

Now define the conjugate momenta

$$p_i = \frac{\partial L}{\partial \dot{q}_i} \quad \text{and these equations become}$$

$$\dot{p}_i = \frac{\partial L}{\partial q_i} \quad (1)$$

$$\text{If } H(q_i, p_i, t) = p_i \dot{q}_i - L(q_i, \dot{q}_i(q_i, p_i, t), t)$$

$$\text{then } \frac{\partial H}{\partial q_i} = - \frac{\partial L}{\partial q_i} + p_j \frac{\partial \dot{q}_j}{\partial q_i} - \frac{\partial L}{\partial \dot{q}_j} \frac{\partial \dot{q}_j}{\partial q_i} = - \frac{\partial L}{\partial q_i} \quad (2)$$

$$\text{Also } \frac{\partial H}{\partial p_i} = \dot{q}_i + p_j \frac{\partial \dot{q}_j}{\partial p_i} - \frac{\partial L}{\partial \dot{q}_j} \frac{\partial \dot{q}_j}{\partial p_i} = \dot{q}_i \quad (3)$$

$$\begin{aligned} (3) &\Rightarrow \dot{q}_i = \frac{\partial H}{\partial p_i} \\ (1) + (2) &\Rightarrow \dot{p}_i = - \frac{\partial H}{\partial q_i} \end{aligned} \quad \left. \vphantom{\begin{aligned} (3) &\Rightarrow \dot{q}_i = \frac{\partial H}{\partial p_i} \\ (1) + (2) &\Rightarrow \dot{p}_i = - \frac{\partial H}{\partial q_i} \end{aligned}} \right\} \text{Hamilton's eqns of motion.}$$

Q1. (b) Show that the following Hamiltonians are integrable and give their integrals of motion.

(i)

$$H(q, p, t) = \frac{p^2}{2} + \cos(q) + e^{-t}$$

(ii)

$$H(\theta_1, \theta_2, I_1, I_2) = I_1^2 + 2I_1I_2 + I_2^2 + I_1I_2 \cos(\theta_1 - 3\theta_2)$$

(7 marks)

i) Since this is a 1 dof system integrability is ensured if there is one constant of the motion.

Here that is  $\bar{H} = \frac{p^2}{2} + \cos(q)$  which is an equally good Hamiltonian for the system since  $e^{-t}$  is a complete differential.

Alternatively use a generating function

$$F = qP + e^{-t}$$

to transform to "new" coordinates  $P = p, Q = q$  with new Hamiltonian  $\bar{H} = H + \frac{\partial F}{\partial t}$ .

ii) Transform to new conjugate variables with  $\phi_1 = \theta_1 - 3\theta_2, \phi_2 = \theta_2$  via.

$$F_2(\theta_i, J_i) = J_1(\theta_1 - 3\theta_2) + J_2\theta_2 \Rightarrow I_1 = J_1, I_2 = -3J_1 + J_2$$

The Hamiltonian

$$H(\phi_i, J_i) = H_0(J_i) + J_1(J_2 - 3J_1) \cos \phi_1$$

is now independent of  $\phi_2 \Rightarrow \frac{\partial H}{\partial \phi_2} = 0 \Rightarrow J_2$  constant of motion.

So, for this 2 dof system, there are 2 constants of the motion:  $H$  and  $J_2 = I_2 + 3I_1$

Question 2 on next page.

$\Rightarrow$  integrable.

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(Actually  $\{J_2, H\} = 0 \Rightarrow$  completely integrable.)

Q2. Give the primary resonance conditions for the following Hamiltonian.

$$H(\theta_1, \theta_2, I_1, I_2) = 2I_1^2 + I_2^2 + \epsilon(I_1^2 \sin(\theta_1)^2 \cos(\theta_2) + I_2^2)$$

If  $H$  is fixed as 1, which tori of the unperturbed system, that is the system with  $\epsilon = 0$ , will be resonant?

Assuming you are not close to resonance give the near identity generating function that removes all the oscillating terms to order  $\epsilon$ . Hence give the approximate integrals of motion in the original variables.

(10 marks)

Since  $\sin^2 \theta_1 \cos \theta_2 = \frac{\cos \theta_2}{2} - \frac{(\cos(2\theta_1 + \theta_2) + \cos(2\theta_1 - \theta_2))}{4}$

There are three resonance conditions (primary)

$$\frac{\partial H_0}{\partial I_2} = 0, \quad \frac{2\partial H_0}{\partial I_1} + \frac{\partial H_0}{\partial I_2} = 0 \quad \text{or} \quad 2\frac{\partial H_0}{\partial I_1} - \frac{\partial H_0}{\partial I_2} = 0$$

$$\Rightarrow I_2 = 0, \quad 4I_1 + 2I_2 = 0 \quad \text{or} \quad 4I_1 - 2I_2 = 0.$$

If  $H_0 = 2I_1^2 + I_2^2 = 1$  then there imply that

$$\left\{ I_1 = \pm \frac{1}{\sqrt{2}}, I_2 = 0 \right\}, \quad \left\{ I_1 = \pm \frac{1}{\sqrt{6}}, I_2 = \pm \frac{2}{\sqrt{6}} \right\}, \quad \left\{ I_1 = \pm \frac{1}{\sqrt{6}}, I_2 = \pm \frac{2}{\sqrt{6}} \right\}$$

Usually the action is taken as  $> 0$

$$\Rightarrow \left\{ I_1 = \frac{1}{\sqrt{2}}, I_2 = 0 \right\}, \text{ No soln }, \left\{ I_1 = \frac{1}{\sqrt{6}}, I_2 = \frac{2}{\sqrt{6}} \right\}.$$

Space for question 2

$$H = 2I_1^2 + I_2^2 + \epsilon \left( I_2^2 + \frac{I_1^2}{4} (2\cos\theta_2 - \cos(2\theta_1 - \theta_2) - \cos(2\theta_1 + \theta_2)) \right)$$

Use a near identity generating function

$$F_2(\theta_i, J_i) = \theta_1 J_1 + \theta_2 J_2$$

$$+ \epsilon (g_1 \sin\theta_2 + g_2 \sin(2\theta_1 - \theta_2) + g_3 \sin(2\theta_1 + \theta_2))$$

to remove all the oscillating terms to order  $\epsilon$ ,

where  $g_i(J_i)$

$$\Rightarrow I_1 = J_1 + \epsilon (g_2 \cos(2\theta_1 - \theta_2) - g_3 \cos(2\theta_1 + \theta_2))$$

$$I_2 = J_2 + \epsilon (g_1 \cos\theta_2 + 2g_2 \cos(2\theta_1 - \theta_2) + 2g_3 \cos(2\theta_1 + \theta_2))$$

To order  $\epsilon$  the Hamiltonian becomes.

$$H = 2J_1^2 + J_2^2 - \epsilon 4J_1 (g_2 \cos(\ ) + g_3 \cos(\ )) + \epsilon 2J_2 (g_1 \cos\theta_2 + 2g_2 \cos(\ ) + 2g_3 \cos(\ )) + \epsilon \left( J_2^2 + \frac{J_1^2}{4} (2\cos\theta_2 - \cos(\ ) - \cos(\ )) \right)$$

$$\text{So choose } 2g_1 J_2 = -\frac{J_1^2}{2} \quad \underline{\underline{g_1 = \frac{J_1^2}{4J_2}}}$$

$$(-4J_1 + 4J_2)g_2 = \frac{J_1^2}{4} \Rightarrow g_2 = \frac{J_1^2}{16(J_2 - J_1)}$$

$$(4(J_1 + J_2))g_3 = \frac{J_1^2}{4} \Rightarrow g_3 = \frac{J_1^2}{16(J_1 + J_2)}$$

The approximate integrals of motion are  $J_1$  and  $J_2$

$$J_1 = I_1 + \epsilon \left( \frac{I_1^2}{16(I_2 - I_1)} \cos(2\phi_1 - \phi_2) - \frac{I_1^2}{16(I_2 + I_1)} \cos(2\phi_1 + \phi_2) \right)$$

Question 3 see next page.

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$$J_2 = I_2 + \epsilon \left( \frac{I_1^2}{4I_2} \cos\phi_2 + \frac{2I_1^2}{16(I_2 - I_1)} \cos(2\phi_1 - \phi_2) + \frac{2I_1^2}{16(I_2 + I_1)} \cos(2\phi_1 + \phi_2) \right)$$

Q3. (a) Describe what happens as you perturb the following twist map

$$I_{n+1} = I_n$$

$$\theta_{n+1} = \theta_n + 2\pi\omega(I_n), \quad \text{mod}(1)$$

if  $\omega = \frac{N}{M}$ , given that  $\omega(I_n)$  is continuous function and the map remains area preserving. (5 marks)

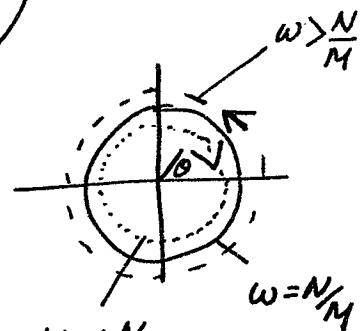
If the unperturbed map is  $T$

$$\begin{pmatrix} \theta_{n+1} \\ I_{n+1} \end{pmatrix} = T \begin{pmatrix} \theta_n \\ I_n \end{pmatrix} = \begin{pmatrix} \theta_n + 2\pi\omega(I_n) \\ I_n \end{pmatrix}$$

Then for  $\omega = \frac{N}{M}$   $T^M$  is the identity.

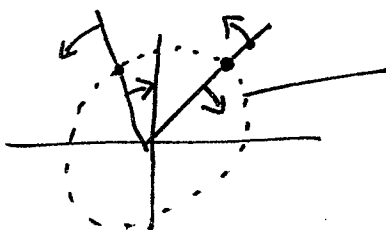
If  $\omega > \frac{N}{M}$  then  $T^M$  twists anticlockwise

But for  $\omega < \frac{N}{M}$   $T^M$  twists clockwise



Now Perturb this to  $T_\epsilon$ . The relative twists will be preserved, i.e. further out the twist is anticlockwise but further in " " " clockwise.

So on each radial line there is a point which is not twisted!



Let the curve of these points be  $R_\epsilon$

Consider  $T_\epsilon^M R_\epsilon$ , each point moves radially.

But area is preserved so

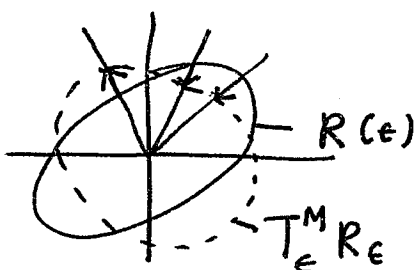
$$T_\epsilon^M R_\epsilon \cap R_\epsilon \neq \emptyset$$

in fact they must intersect in an even no of points.

Further if  $x$  is an intersection point then so is  $T_\epsilon x \Rightarrow \exists$  at least  $2kM$

~~period~~ fixed points of  $T_\epsilon^M$  or period- $M$

points of  $T_\epsilon$ .



Q3. (b) Show that the following map

$$u_{n+1} = u_n - A \sin(2\pi\phi_n)$$

$$\phi_{n+1} = \phi_n + (u_{n+1})^2 \pmod{1}$$

is a product of two involutions and find any lines or curves which are fixed under the involutions of the map. Using this or otherwise investigate the possibility of symmetric period-2 points where one iterate lies on the line  $\phi_n = 0$ .

(5 marks)

First find the involutions:

$$I_1 \begin{pmatrix} \phi \\ u \end{pmatrix} = \begin{pmatrix} -\phi \\ u - A \sin 2\pi\phi \end{pmatrix} \text{ then } I_1^2 \begin{pmatrix} \phi \\ u \end{pmatrix} = \begin{pmatrix} -(-\phi) \\ u - A \sin 2\pi(-\phi) - A \sin 2\pi\phi \end{pmatrix} = \begin{pmatrix} \phi \\ u \end{pmatrix}$$

$$I_2 \begin{pmatrix} \phi \\ u \end{pmatrix} = \begin{pmatrix} -\phi + u^2 \\ u \end{pmatrix} \text{ then } I_2^2 \begin{pmatrix} \phi \\ u \end{pmatrix} = \begin{pmatrix} -(-\phi + u^2) + u^2 \\ u \end{pmatrix} = \begin{pmatrix} \phi \\ u \end{pmatrix}$$

such that  $I_2 I_1$  is the map:

$$I_2 I_1 \begin{pmatrix} \phi \\ u \end{pmatrix} = \begin{pmatrix} -(-\phi) + (u - A \sin 2\pi\phi)^2 \\ u - A \sin 2\pi\phi \end{pmatrix} = \begin{pmatrix} \phi + (u - A \sin 2\pi\phi)^2 \\ u - A \sin 2\pi\phi \end{pmatrix}$$

Fixed points of  $I_1$

$$\phi = -\phi \text{ and } u = u - A \sin 2\pi\phi \Rightarrow \phi = 0 \text{ or } \frac{1}{2} \text{ are fixed lines}$$

Fixed points of  $I_2$

$$\phi = -\phi + u^2 \Rightarrow u^2 = 2\phi \pmod{1} \Rightarrow \left. \begin{aligned} \phi &= \frac{u^2}{2} \\ \phi &= \frac{1}{2} + \frac{u^2}{2} \end{aligned} \right\} \text{ fixed curves.}$$

There is a symmetric period-2 point at

$$u_n = \frac{1}{\sqrt{2}}, \phi_n = 0 \Rightarrow u_{n+1} = \frac{1}{\sqrt{2}}, \phi_{n+1} = \frac{1}{2}$$

$$u_{n+2} = \frac{1}{\sqrt{2}}, \phi_{n+2} = 1/\text{mod}(1) = 0.$$

There are various ways to get this; Question 4 see next page. TURN OVER

directly by considering period-2 pts with  $\phi_n = 0$ ,  
or by using properties of involutions.